

# Canonical Mapping of $O(N)$ Vector /Higher Spin Correspondence



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# Higher Spin Gravity

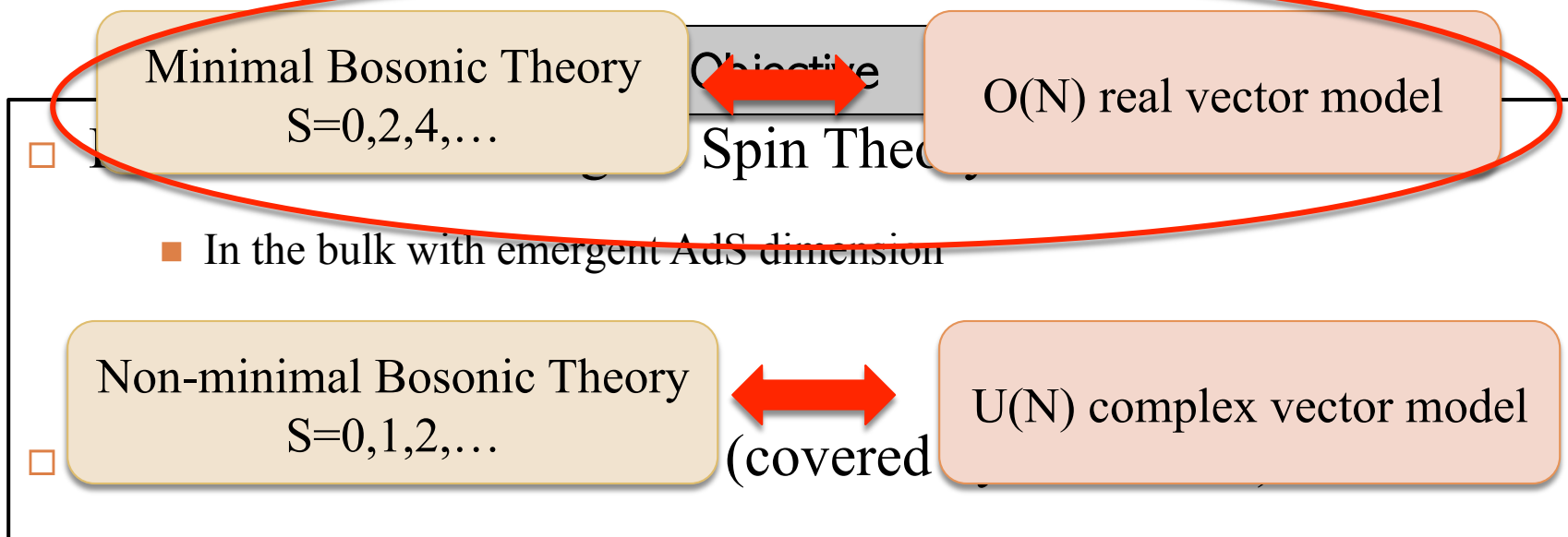
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## AdS<sub>4</sub> Higher Spin Theory

- M. A. Vasiliev
- Infinite tower of higher spin

## CFT<sub>3</sub> : O(N) , U(N) vector

- Boson :  $\varphi_i$  ( $i=1,\dots,N$ )
- Singlet sector





# Contents

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- 1. Review : Map in Light-cone Gauge

[A. Jevicki, J.P.Rodrigues, R.d.M.Koch, K.Jin, 2010]

- 2. Extention to Time-like Gauge

[JY, A. Jevicki, J.P.Rodrigues, R.d.M.Koch, 2014]

- Quick Derivation of Bi-local Map
- $SO(2,3)$  generator for HS field and Bi-local Map



# HS Field : Light-cone Gauge

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- Spin-s field : package with harmonic oscillator

$$\Phi = \Phi^{A_1 \cdots A_s} \alpha_{A_1} \cdots \alpha_{A_s} \quad \Phi^{A_1 \cdots A_s} = e_{\mu_1}^{A_1} \cdots e_{\mu_s}^{A_s} h^{\mu_1 \cdots \mu_s} \quad [\bar{\alpha}_A, \alpha_B] = \eta_{AB}$$

- Symmetric, double-traceless

$$(\bar{\alpha}^A \bar{\alpha}_A)^2 \Phi = 0 \quad \alpha^A \bar{\alpha}_A \Phi = s \Phi$$

- Light-cone gauge :  $\bar{\alpha}^+ \Phi = 0$

- Physical Field : two helicities

$$\Phi_{\text{phy}} = \Phi_{\text{phy}}^{I_1 \cdots I_s} \alpha_{I_1} \cdots \alpha_{I_s} \quad (I = 1, 2) \quad [\text{R.R.Metsaev, 1999}]$$

$$\bar{\alpha}^I \bar{\alpha}_I \Phi_{\text{phy}} = 0 \quad \alpha^A \bar{\alpha}_A \Phi_{\text{phy}} = s \Phi_{\text{phy}} \quad \partial^\mu \partial_\mu \Phi_{\text{phy}} = 0$$

$$\text{Spin matrix } M_{12} = \alpha_1 \bar{\alpha}_2 - \alpha_2 \bar{\alpha}_1 \longrightarrow S^1$$

$$\Phi_{\text{phy}} = \Phi_s \beta_+^s + \bar{\Phi}_s \beta_-^s \quad \beta_\pm = \alpha_1 \pm i \alpha_2$$





# HS SO(2,3) Generators : Light-cone

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- SO(2,3) generator for physical HS field on  $\text{AdS}_4 \times S^1$  (in Light-cone gauge) [R.R.Metsaev, 1999]

Poincare

$$\begin{aligned}\hat{p}^- &= -\frac{p^x p^x + p^z p^z}{2p^+} & \hat{m}^{+-} &= -x^- p^+ \\ \hat{p}^+ &= p^+ & \hat{m}^{+x} &= -x p^+ \\ \hat{p}^x &= p^x & \hat{m}^{-x} &= x^- p^x - x \hat{p}^- + \frac{p^\theta p^z}{p^+}\end{aligned}$$

Dilation

$$\hat{d} = x^- p^+ + x p^x + z p^z + 1$$

Special  
Conformal

$$\begin{aligned}\hat{k}^- &= -\frac{1}{2}(x^2 + z^2)\hat{p}^- + x^- \hat{d} + \frac{1}{p^+}((x p^z - z p^x)p^\theta + (p^\theta)^2) \\ \hat{k}^x &= -\frac{1}{2}(x^2 + z^2)p^x + x \hat{d} + z p^\theta \\ \hat{k}^+ &= -\frac{1}{2}(x^2 + z^2)p^+\end{aligned}$$



# Schematic Bi-local Map

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- SO(2,3) Decomposition [C. Fronsdal & M. Flato, 1978]

$$D(1/2, 0) \otimes D(1/2, 0) = \sum_{s=0}^{\infty} D(s+1, s) \text{ cf. Addition of Angular momentum}$$

- Bi-local Field : O(N) invariants

$$\Psi(x^+; x_1^-, x_1, x_2^-, x_2) = \vec{\varphi}(x^+; x_1^-, x_1) \cdot \vec{\varphi}(x^+; x_2^-, x_2)$$

- SO(2,3) Generators for Bi-local field

- Map from collective field to HS field

$$H(X, y) = \int du_1 du_2 K(X, y|u_1, u_2) \Psi(u_1, u_2)$$



# Bi-local Generator : Light-cone Gauge

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## □ SO(2,3) Bi-local generator

Poincare

$$\hat{p}^- = -\frac{p_1 p_1}{2p_1^+} - \frac{p_2 p_2}{2p_2^+}$$

$$\hat{p}^+ = p_1^+ + p_2^+$$

$$\hat{p}^x = p_1 + p_2$$

$$\hat{m}^{+x} = -x_1 p_1^+ - x_2 p_2^+$$

$$\hat{m}^{-x} = x_1^- p_1 + x_2^- p_2 - x_1 \left( -\frac{p_1 p_1}{2p_1^+} \right) - x_2 \left( -\frac{p_2 p_2}{2p_2^+} \right)$$

$$\hat{m}^{+-} = -x_1^- p_1^+ - x_2^- p_2^+$$

Dilation

$$\hat{d} = x_1^- p_1^+ + x_1 p_1 + x_2^- p_2^+ + x_2 p_2 + 1$$

Special  
Conformal

$$\hat{k}^+ = -\frac{1}{2} x_1 x_1 p_1^+ - \frac{1}{2} x_2 x_2 p_2^+$$

$$\hat{k} = -\frac{1}{2} x_1 x_1 p_1 + x_1 (x_1^- p_1^+ + x_1 p_1 + \frac{1}{2}) - \frac{1}{2} x_2 x_2 p_2 + x_2 (x_2^- p_2^+ + x_2 p_2 + \frac{1}{2})$$

$$\hat{k}^- = -\frac{1}{2} x_1 x_1 \left( -\frac{p_1 p_1}{2p_1^+} \right) + x_1^- (x_1^- p_1^+ + x_1 p_1 + \frac{1}{2}) - \frac{1}{2} x_1 x_1 \left( -\frac{p_1 p_1}{2p_1^+} \right) + x_2^- (x_2^- p_2^+ + x_2 p_2 + \frac{1}{2})$$



# Map : Canonical Transformation

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- Identifying two generator  $J_{\text{HS}} = J_{\text{Bi-local}} = J_{\text{CFT}}^1 + J_{\text{CFT}}^2$
- Bi-local Map [A. Jevicki, J.P.Rodrigues, R.d.M.Koch, K.Jin, 2010]

$$\begin{aligned}
 p^+ &= p_1^+ + p_2^+ \\
 p &= p_1 + p_2 \\
 \text{AdS } p^z &= p_1 \sqrt{\frac{p_2^+}{p_1^+}} - p_2 \sqrt{\frac{p_1^+}{p_2^+}} \quad \text{CFT} \\
 \theta &= 2 \arctan \sqrt{\frac{p_2^+}{p_1^+}}
 \end{aligned}$$

$$\begin{aligned}
 x^- &= \frac{x_1^- p_1^+ + x_2^- p_2^+}{p_1^+ + p_2^+} \\
 x &= \frac{x_1 p_1^+ + x_2 p_2^+}{p_1^+ + p_2^+} \\
 \text{AdS } z &= \frac{(x_1 - x_2) \sqrt{p_1^+ p_2^+}}{p_1^+ + p_2^+} \quad \text{CFT} \\
 p^\theta &= \sqrt{p_1^+ p_2^+} (x_1^- - x_2^-) + \frac{x_1 - x_2}{2} \left( p_1 \sqrt{\frac{p_2^+}{p_1^+}} + p_2 \sqrt{\frac{p_1^+}{p_2^+}} \right)
 \end{aligned}$$

- Canonical Transformation

$$\{p^+, x^-\}_{\text{Bi-local PB}} = \{p^x, x\}_{\text{Bi-local PB}} = \{p^z, z\}_{\text{Bi-local PB}} = \{p^\theta, \theta\}_{\text{Bi-local PB}} = 1$$



# Kernel

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□ Kernel :  $\mathcal{K}(p^+, p, p^z, \theta; p_1^+, p_1, p_2^+, p_2) = \mathcal{J} \delta(p_1^+ + p_2^+ - p^+) \delta(p_1 + p_2 - p)$   
 $\times \delta(p_1 \sqrt{\frac{p_2^+}{p_1^+}} - p_2 \sqrt{\frac{p_1^+}{p_2^+}} - p^z) \delta(2 \arctan \sqrt{\frac{p_2^+}{p_1^+}} - \theta)$

□ One-to-one, induce Extra Bulk coordinate  $z$

□ Transformation : From Bi-local field to HS field

$$\tilde{\mathcal{H}}(p^+, p, p^z, \theta) = \int dp_1^+ dp_1 dp_2^+ dp_2 \mathcal{K}(p^+, p, p^z, \theta; p_1^+, p_1, p_2^+, p_2) \tilde{\Psi}(p_1^+, p_1, p_2^+, p_2)$$

□ Local

$$\tilde{\mathcal{H}}(p^+, p, p^z, \theta) = \int dp_1^+ dp_1 dp_2^+ dp_2 \mathcal{K} \tilde{\Psi}$$

$$\tilde{\mathcal{W}}(p^+, p, p^z, \theta) = \int dp_1^+ dp_1 dp_2^+ dp_2 \mathcal{K} \tilde{\Pi}$$

$$[\tilde{\Psi}(p_1^+, p_1, p_2^+, p_2), \tilde{\Pi}(k_1^+, k_1, k_2^+, k_2)] = p_1^+ p_2^+ \delta(p_1^+ - k_1^+) \delta(p_1 - k_1) \delta(p_2^+ - k_2^+) \delta(p_2 - k_2)$$



$$[\tilde{\mathcal{H}}(p^+, p, p^z, \theta), \tilde{\mathcal{W}}(k^+, k, k^z, \phi)] = p^+ \delta(p^+ - k^+) \delta(p - k) \delta(p^z - k^z) \delta(\theta - \phi)$$



# Spin-s Primary Current

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- Fourier Transformation :  $(p^+, p, p^z, \theta) \longrightarrow (x^-, x, z, p^\theta = s)$

$$\mathcal{H}_s(x^+, x^-, x, z) = \int_{-2p^+p^- - p^2 > 0} dp^- dp^+ dp e^{ix^+p^- + ix^-p^+ + ixp} J_{-\frac{1}{2}} \left( z \sqrt{-2p^+p^- - p^2} \right) \\ \times \sqrt{\frac{\pi z \sqrt{-2p^+p^- - p^2}}{2}} \frac{s!}{\Gamma(s + \frac{1}{2}) (p^+)^s} \tilde{\mathcal{O}}_s(p^-, p^+, p) \quad [\text{I.Bena, 1999}]$$

$$\xrightarrow{z \longrightarrow 0} \frac{s!}{\Gamma(s + \frac{1}{2})} \partial_-^s \int_{-2p^+p^- - p^2 > 0} dp^- dp^+ dp e^{ix^+p^- + ix^-p^+ + ixp} \tilde{\mathcal{O}}_s(x^-, x^+, x)$$

$$\mathcal{O}_s = \sum_{n=0}^{s/2} \frac{(-4)^n}{(2n)!} \sum_{k=0}^{s-2n} \frac{(-1)^k}{k!(s-2n-k)!} \partial_-^{n+k} \vec{\phi} \cdot \partial_-^{s-n-k} \vec{\phi} : \text{Spin-s primary field in CFT}$$

- Bi-local map (in  $z \longrightarrow 0$  limit) gives primary CFT operator



# Time-like Gauge : Quick Derivation

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- Transformation in “Momentum” space

$$(\vec{p}_1, \vec{p}_2) \longrightarrow (\vec{p}, p^z, \theta)$$

- Total momentum ( CM of two particles )

$$\vec{p} = \vec{p}_1 + \vec{p}_2 \quad p^0 = |\vec{p}_1| + |\vec{p}_2|$$

- On-shell condition

$$p^z = \pm \sqrt{(p^0)^2 - \vec{p}^2} = 2\sqrt{|\vec{p}_1||\vec{p}_2|} \sin\left(\frac{\varphi_1 - \varphi_2}{2}\right)$$

- SO(2,3) Casimir

$$C_{SO(2,3)} = E_0^2 + s^2 \xrightarrow{\text{Massless representation } E_0 = s + 1} C_{SO(2,3)} = 2s^2 = 2(p^\theta)^2$$

- Compare to SO(2,3) Casimir of Bi-local generator

$$p^\theta = \sqrt{|\vec{p}_1||\vec{p}_2|} \cos \frac{\varphi_1 + \varphi_2}{2} (x_2^1 - x_1^1) + \sqrt{|\vec{p}_1||\vec{p}_2|} \sin \frac{\varphi_1 + \varphi_2}{2} (x_2^2 - x_1^2)$$



# Time-like Gauge : Quick Derivation

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- Ansatz for  $\theta$  ( to have canonical transformation, )

$$\theta = \arctan \left( \frac{2\vec{p}_2 \times \vec{p}_1}{(|\vec{p}_2| - |\vec{p}_1|)p^z} \right)$$

- Can be found naturally.(Later)

- Bi-local Map in time-like gauge

$$\begin{aligned} \mathcal{K}(\vec{p}, p^z, \theta; \vec{p}_1, \vec{p}_2) = & \mathcal{J}(\vec{p}_1, \vec{p}_2) \delta^{(2)}(\vec{p}_1 + \vec{p}_2 - \vec{p}) \delta \left( 2\sqrt{|\vec{p}_1| |\vec{p}_2|} \sin \left( \frac{\varphi_1 - \varphi_2}{2} \right) - p^z \right) \\ & \times \delta \left( \arctan \left( \frac{2\vec{p}_2 \times \vec{p}_1}{(|\vec{p}_2| - |\vec{p}_1|)p^z} \right) - \theta \right) \end{aligned}$$





# Coordinates Transformation

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## □ “Momentum” space transformation

→ Coordinates transformation by Chain rule

$$x^1 = \frac{|\vec{p}_1| x_1^1 + |\vec{p}_2| x_2^1}{\sqrt{(p^z)^2 + \vec{p}^2}} - \frac{p^2 p^z p^\theta}{\vec{p}^2 \sqrt{(p^z)^2 + \vec{p}^2}}$$

$$x^2 = \frac{|\vec{p}_1| x_1^2 + |\vec{p}_2| x_2^2}{\sqrt{(p^z)^2 + \vec{p}^2}} + \frac{p^1 p^z p^\theta}{\vec{p}^2 \sqrt{(p^z)^2 + \vec{p}^2}}$$

$$z = \frac{(\vec{x}_1 - \vec{x}_2) \cdot \vec{p}_1 |\vec{p}_2| - (\vec{x}_1 - \vec{x}_2) \cdot \vec{p}_2 |\vec{p}_1|}{p^z (|\vec{p}_1| + |\vec{p}_2|)}$$

$$p^\theta = \sqrt{|\vec{p}_1| |\vec{p}_2|} \cos \frac{\varphi_1 + \varphi_2}{2} (x_2^1 - x_1^1) + \sqrt{|\vec{p}_1| |\vec{p}_2|} \sin \frac{\varphi_1 + \varphi_2}{2} (x_2^2 - x_1^2)$$

## □ Canonical Transformation by Construction

$$\{p^i, x^j\} = \delta^{ij} \quad \{p^z, z\} = \{p^\theta, \theta\} = 1$$



# Reduction

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- Package spin  $s$  current with  $\alpha$

$$\mathcal{O}_s(x; \alpha) = \mathcal{O}_{\mu_1 \mu_2 \dots \mu_s}^s \alpha^{\mu_1} \dots \alpha^{\mu_s} \quad [\bar{\alpha}^\mu, \alpha^\nu] = \eta^{\mu\nu}$$

- Constraints

$$\bar{\alpha}^\mu \partial_\mu \mathcal{O}_s(x; \alpha) = 0 \quad : \text{Conservation}$$

$$\bar{\alpha}^\mu \bar{\alpha}_\mu \mathcal{O}_s(x; \alpha) = 0 \quad : \text{Traceless}$$

$$\alpha^\mu \bar{\alpha}_\mu \mathcal{O}_s(x; \alpha) = s \mathcal{O}^s(x; \alpha) \quad : \text{Spin } s$$

- To make contact with physical collective picture, we need to solve constraints
- Let's solve



# Solve Constraints : Canonical Transformation

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□ (6+6) dim phase space  $(x^\mu, p^\mu; \alpha^\mu, \bar{\alpha}^\mu)$

□ Canonical Transformation

$$\bar{\alpha}^\mu \partial_\mu \mathcal{O}_s(x; \alpha) = 0$$

$$\bar{\alpha}^\mu \bar{\alpha}_\mu \mathcal{O}_s(x; \alpha) = 0$$



$$p^\mu \bar{\alpha}_\mu = 0$$

$$\bar{\alpha}^\mu \bar{\alpha}_\mu = 0$$



$T_1$

$$(\bar{\beta}^1)^2 + (\bar{\beta}^2)^2 - \frac{1}{(\bar{p}^0)^2} (\bar{\beta}^0 \bar{p}^0 + \bar{\beta}^1 \bar{p}^1 + \bar{\beta}^2 \bar{p}^2)^2 = 0$$

$$-\bar{\beta}^0 \bar{p}^0 = 0$$



$T_2$

$$2\bar{\gamma}^1 \bar{\gamma}^2 - \bar{\gamma}^0 f(\tilde{p}, \tilde{\gamma}) = 0$$

$$-\bar{\gamma}^0 \tilde{p}^0 = 0$$

$$(x^\mu, p^\mu; \alpha^\mu, \bar{\alpha}^\mu)$$



$T_1$

$$(\bar{x}^\mu, \bar{p}^\mu; \beta^\mu, \bar{\beta}^\mu)$$



$T_2$

$$(\tilde{x}^\mu, \tilde{p}^\mu; \gamma^\mu, \bar{\gamma}^\mu)$$

Two  
Solutions

$$\gamma^0 = \bar{\gamma}^0 = \gamma^2 = \bar{\gamma}^2 = 0$$

$$\gamma^0 = \bar{\gamma}^0 = \gamma^1 = \bar{\gamma}^1 = 0$$



# (Unconstrained) Primary spin-s Current

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## □ (Unconstrained) Primary spin-s operator ( $p^0 > 0$ )

$$\begin{aligned} \tilde{\mathcal{O}}_s(p; \epsilon) = & \int d\vec{p}_1 d\vec{p}_2 \delta^{(3)}(p_1^\mu + p_2^\mu - p^\mu) \frac{(-1)^s 2^{\frac{3}{2}s-1}}{s!} \\ & \times (2 |\vec{p}_1| |\vec{p}_2| - 2\vec{p}_1 \cdot \vec{p}_2)^{\frac{s}{2}} e^{is\Theta(\vec{p}_1, \vec{p}_2)} a(p_1, p_2) \end{aligned}$$

: annihilation operator of bi-local field

$$\begin{aligned} \tilde{\mathcal{O}}_s(p; \epsilon^*) = & \int d\vec{p}_1 d\vec{p}_2 \delta^{(3)}(p_1^\mu + p_2^\mu - p^\mu) \frac{(-1)^s 2^{\frac{3}{2}s-1}}{s!} \\ & \times (2 |\vec{p}_1| |\vec{p}_2| - 2\vec{p}_1 \cdot \vec{p}_2)^{\frac{s}{2}} e^{-is\Theta(\vec{p}_1, \vec{p}_2)} a(p_1, p_2) \end{aligned}$$

$$\Theta(\vec{p}_1, \vec{p}_2) \equiv 2 \arctan \left( \frac{2\vec{p}_1 \times \vec{p}_2}{\sqrt{2\sqrt{\vec{p}_1^2 \vec{p}_2^2} - 2\vec{p}_1 \cdot \vec{p}_2} (\sqrt{\vec{p}_1^2} - \sqrt{\vec{p}_2^2})} \right)$$

: Equal to  $\theta$  map (up to sign)

$$\epsilon(p) \equiv \frac{1}{\sqrt{2} |\vec{p}|} \left( \frac{\vec{p}^2}{\sqrt{-p^\mu p_\mu}}, \frac{p^0 p^1}{\sqrt{-p^\mu p_\mu}} + ip^2, \frac{p^0 p^2}{\sqrt{-p^\mu p_\mu}} - ip^1 \right) \quad : \text{polarization vector}$$



# SO(2,3) Generator for HS field

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## □ SO(2,3) Generator for physical spin-s primary

$$\hat{p}^\mu = p^\mu$$

$$\hat{d} = -x^0 p^0 + x^1 p^1 + x^2 p^2 + 2\alpha^1 \bar{\alpha}^1$$

$$\hat{m}^{12} = x^1 p^2 - x^2 p^1$$

$$\hat{m}^{20} = x^2 p^0 - x^0 p^2 + \left( \frac{p^2}{p^0} + i \frac{p^1}{q} - i \frac{(p^0)^2 p^1}{q \bar{p}^2} \right) \alpha^1 \bar{\alpha}^1$$

$$\hat{k}^0 = -\frac{1}{2}(x^\mu x_\mu) p^0 + x^0 \hat{D} + \frac{1}{2}(\alpha^1 \bar{\alpha}^1)^2 \left( \frac{q^2}{p^0 \bar{p}^2} - \frac{3}{p^0} - \frac{\bar{p}^2}{p^0 q^2} \right) - (\alpha^1 \bar{\alpha}^1) \left( \frac{x^1 p^1 + x^2 p^2}{p^0} + i \frac{q}{\bar{p}^2} (x^1 p^2 - x^2 p^1) \right)$$

$$(p^0, p^1, p^2) \Rightarrow (p^1, p^2, p^z = \sqrt{(p^0)^2 - \bar{p}^2})$$

$$\hat{p}^0 = \sqrt{\bar{p}^2 + (p^z)^2}$$

$$\hat{p}^1 = p^1$$

$$\hat{p}^2 = p^2$$

$$\hat{m}^{01} = -x^1 \hat{p}^0 - \frac{p^2 p^z}{\bar{p}^2} p^\theta$$

$$\hat{m}^{12} = x^1 p^2 - x^2 p^1$$

$$\hat{m}^{20} = x^2 \hat{p}^0 - \frac{p^1 p^z}{\bar{p}^2} p^\theta$$

$$\hat{d} = x^1 p^1 + x^2 p^2 + z p^z$$

$$\hat{k}^0 = -\frac{1}{2}(\vec{x}^2 + z^2) \hat{p}^0 - \frac{p^z}{\bar{p}^2} \hat{m}^{12} p^\theta - \frac{\hat{p}^0}{2\bar{p}^2} (p^\theta)^2$$

$$\hat{k}^1 = -\frac{1}{2}(\vec{x}^2 + z^2) p^1 + x^1 D + \frac{z p^2 \hat{p}^0}{\bar{p}^2} p^\theta + \frac{p^1}{2\bar{p}^2} (p^\theta)^2$$

$$\hat{k}^2 = -\frac{1}{2}(\vec{x}^2 + z^2) p^2 + x^2 D - \frac{z p^1 \hat{p}^0}{\bar{p}^2} p^\theta + \frac{p^2}{2\bar{p}^2} (p^\theta)^2$$

HS Generator



# Bi-local Map in Time-like Gauge

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- HS generator = Bi-local generator
- Same Result

$$\vec{p} = \vec{p}_1 + \vec{p}_2 \quad p^z = 2\sqrt{|\vec{p}_1||\vec{p}_2|} \sin\left(\frac{\varphi_1 - \varphi_2}{2}\right)$$

$$p^\theta = \sqrt{|\vec{p}_1||\vec{p}_2|} \cos\frac{\varphi_1 + \varphi_2}{2} (x_2^1 - x_1^1)$$

$$x^1 = \frac{|\vec{p}_1| x_1^1 + |\vec{p}_2| x_2^1}{\sqrt{(p^z)^2 + \vec{p}^2}} - \frac{p^2 p^z p^\theta}{\vec{p}^2 \sqrt{(p^z)^2 + \vec{p}^2}}$$

$$x^2 = \frac{|\vec{p}_1| x_1^2 + |\vec{p}_2| x_2^2}{\sqrt{(p^z)^2 + \vec{p}^2}} + \frac{p^1 p^z p^\theta}{\vec{p}^2 \sqrt{(p^z)^2 + \vec{p}^2}}$$

$$z = \frac{(\vec{x}_1 - \vec{x}_2) \cdot \vec{p}_1 |\vec{p}_2| - (\vec{x}_1 - \vec{x}_2) \cdot \vec{p}_2 |\vec{p}_1|}{p^z (|\vec{p}_1| + |\vec{p}_2|)}$$

$$p^\theta = \sqrt{|\vec{p}_1||\vec{p}_2|} \cos\frac{\varphi_1 + \varphi_2}{2} (x_2^1 - x_1^1) + \sqrt{|\vec{p}_1||\vec{p}_2|} \sin\frac{\varphi_1 + \varphi_2}{2} (x_2^2 - x_1^2)$$

$$+ \sqrt{|\vec{p}_1||\vec{p}_2|} \sin\frac{\varphi_1 + \varphi_2}{2} (x_2^2 - x_1^2)$$



# Properties of Kernel

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## □ “Momentum” space transformation : Kernel

$$\begin{aligned}\mathcal{K}(\vec{p}, p^z, \theta; \vec{p}_1, \vec{p}_2) &= \mathcal{J}(\vec{p}_1, \vec{p}_2) \delta^{(2)}(\vec{p}_1 + \vec{p}_2 - \vec{p}) \delta \left( 2\sqrt{|\vec{p}_1| |\vec{p}_2|} \sin \left( \frac{\varphi_1 - \varphi_2}{2} \right) - p^z \right) \\ &\quad \times \delta \left( \arctan \left( \frac{2\vec{p}_2 \times \vec{p}_1}{(|\vec{p}_2| - |\vec{p}_1|)p^z} \right) - \theta \right) \\ &= \delta^{(2)}(\vec{p}_1 - \vec{\beta}_1(\vec{p}, p^z, \theta)) \delta^{(2)}(\vec{p}_2 - \vec{\beta}_2(\vec{p}, p^z, \theta))\end{aligned}$$

### ■ One-to-one

## □ Inverse kernel

$$\begin{aligned}\mathcal{Q}(\vec{p}_1, \vec{p}_2; \vec{p}, p^z, \theta) &= \delta^{(2)}(\vec{p}_1 + \vec{p}_2 - \vec{p}) \delta \left( 2\sqrt{|\vec{p}_1| |\vec{p}_2|} \sin \left( \frac{\varphi_1 - \varphi_2}{2} \right) - p^z \right) \\ &\quad \times \delta \left( \arctan \left( \frac{2\vec{p}_2 \times \vec{p}_1}{(|\vec{p}_2| - |\vec{p}_1|)p^z} \right) - \theta \right)\end{aligned}$$



# Properties of Kernel

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□ **Local**

$$\tilde{\mathcal{H}}(\vec{p}, p^z, \theta) = \int d\vec{p}_1 d\vec{p}_2 \mathcal{K}(\vec{p}, p^z, \theta; \vec{p}_1, \vec{p}_2) \tilde{\Psi}(\vec{p}_1, \vec{p}_2)$$

$$\tilde{\mathcal{W}}(\vec{p}, p^z, \theta) = \int d\vec{p}_1 d\vec{p}_2 \mathcal{K}(\vec{p}, p^z, \theta; \vec{p}_1, \vec{p}_2) \tilde{\Pi}(\vec{p}_1, \vec{p}_2)$$

$$[\tilde{\Psi}(\vec{p}_1, \vec{p}_2), \tilde{\Pi}(\vec{k}_1, \vec{k}_2)] = |\vec{p}_1| |\vec{p}_2| \delta^{(2)}(\vec{p}_1 - \vec{k}_1) \delta^{(2)}(\vec{p}_2 - \vec{k}_2)$$



$$[\tilde{\mathcal{H}}(\vec{p}, p^z, \theta), \tilde{\mathcal{W}}(\vec{k}, k^z, \phi)] = \sqrt{\vec{p}^2 + (p^z)^2} \delta^{(2)}(\vec{p} - \vec{k}) \delta(p^z - k^z) \delta(\theta - \phi)$$

□ **Redundancy in Bi-local field**  $\tilde{\Psi}(\vec{p}_1, \vec{p}_2) = \tilde{\Psi}(\vec{p}_2, \vec{p}_1)$

→ **Redundancy in HS field**  $\tilde{\mathcal{H}}(\vec{p}, p^z, \theta) = \tilde{\mathcal{H}}(\vec{p}, -p^z, -\theta)$

$$\cos s\theta \cos p^z z \quad \sin s\theta \sin p^z z \quad \sin s\theta \cos p^z z \quad \cos s\theta \sin p^z z$$

$$\mathcal{H}_s^{(+)}(t, \vec{x}, z) \sim \cos s\theta \cos p^z z$$

$$\mathcal{H}_s^{(-)}(t, \vec{x}, z) \sim \sin s\theta \sin p^z z \quad \mathcal{H}_0^{(-)} = 0$$





# From CFT Primary to HS Field

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- Kernel : creation(annihilation) operator of bi-local field to HS creation(annihilation) operator
- Quadratic Hamiltonian of Collective field is mapped to Quadratic Hamiltonian of HS field

$$H_2 = \frac{1}{2} \int_{\vec{p}_1 \geq \vec{p}_2} d\vec{p}_1 d\vec{p}_2 \left[ \tilde{\Pi}'(\vec{p}_1, -\vec{p}_2) \tilde{\Pi}'(\vec{p}_2, -\vec{p}_1) + \omega_{\vec{p}_1, \vec{p}_2}^2 \tilde{\Psi}'(\vec{p}_1, -\vec{p}_2) \tilde{\Psi}'(\vec{p}_2, -\vec{p}_1) \right]$$



$$H_2 = \int d\vec{p} dp^z d\theta \sqrt{\vec{p}^2 + (p^z)^2} \mathcal{A}^\dagger(\vec{p}, p^z, \theta) \mathcal{A}(\vec{p}, p^z, \theta)$$



# From CFT Primary to HS Field

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- Fourier Transformation :  $(\vec{p}, p^z, \theta) \longrightarrow (\vec{x}, z, p^\theta = s)$
- New basis

$$\mathcal{H}_s^{(\pm)}(\vec{x}, z) \equiv \frac{1}{2} [\mathcal{H}_s(\vec{x}, z) \pm \mathcal{H}_{-s}(\vec{x}, z)]$$

- Perform time-evolution and change of variable
  - Recall time-evolution of HS field is induced from time-evolution of bi-local field
  - From  $p^z$  to  $p^0 = \sqrt{\vec{p}^2 + (p^z)^2}$

$$\mathcal{H}_s^{(\pm)}(t; \vec{x}, z) = \int d\vec{p} dp^z \dots = \int d\vec{p} dp^0 \dots$$



# Summary : Time-like Map

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## □ Canonical Map to Bulk field

$$\mathcal{H}_s^{(+)}(t; \vec{x}, z) = \frac{\sqrt{\pi} s!}{2^{\frac{3}{2}s + \frac{3}{2}}} \int_{p^0 > |\vec{p}|} \frac{d^2 \vec{p} dp^0}{(2\pi)^3 2p^0} e^{-ip^0 t + i\vec{p} \cdot \vec{x}} z^{\frac{1}{2}} [(p^0)^2 - \vec{p}^2]^{-\frac{s}{2} + \frac{1}{4}} \\ \times J_{-\frac{1}{2}}(\sqrt{(p^0)^2 - \vec{p}^2} z) \left[ \tilde{\mathcal{O}}_s(p; \epsilon) + \tilde{\mathcal{O}}_s(p; \epsilon^*) \right] + \text{h.c}$$

$$\xrightarrow{z \rightarrow 0} 2^{-\frac{3}{2}s-1} s! \left( -(\partial^0)^2 + \vec{\partial}^2 \right)^{-\frac{s}{2}} \int_{p^0 > |\vec{p}|} \frac{d^2 \vec{p} dp^0}{(2\pi)^3 2p^0} \left[ \tilde{\mathcal{O}}_s(p; \epsilon) + \tilde{\mathcal{O}}_s(p; \epsilon^*) \right] + \text{h.c.}$$

$$\tilde{\mathcal{O}}_s(p; \epsilon) = \int d\vec{p}_1 d\vec{p}_2 \delta^{(3)}(p_1^\mu + p_2^\mu - p^\mu) \frac{(-1)^s 2^{\frac{3}{2}s-1}}{s!} \\ \times (2 |\vec{p}_1| |\vec{p}_2| - 2\vec{p}_1 \cdot \vec{p}_2)^{\frac{s}{2}} e^{is\Theta(\vec{p}_1, \vec{p}_2)} a(p_1, p_2)$$

: Physical spin-s primary in CFT



# Solving Constraints

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□ Define Kernel  $\mathcal{M}_{tot} \equiv \mathcal{M}_1 \mathcal{M}_2 \mathcal{M}_3 \mathcal{M}_4$   $\tan \theta = \frac{p^2}{p^1}$

$$\mathcal{M}_1 \equiv \exp \left[ -\alpha^0 \left( \frac{1}{p^0} (\bar{\alpha}^1 p^1 + \bar{\alpha}^2 p^2) \right) \right] \quad \mathcal{M}_3 \equiv \exp \left[ -\alpha^1 \bar{\alpha}^1 \log \left( \frac{q}{p^0} \right) \right]$$

$$\mathcal{M}_2 \equiv \exp \left[ -\theta (\alpha^1 \bar{\alpha}^2 - \alpha^2 \bar{\alpha}^1) \right] \quad \mathcal{M}_4 \equiv \exp \left[ \frac{\pi}{2} i \alpha^2 \bar{\alpha}^2 \right] \exp \left[ \frac{\pi}{4} (\alpha^1 \bar{\alpha}^2 - \alpha^2 \bar{\alpha}^1) \right]$$

■ Designed to generate previous canonical transformation

□ Spin-s current :  $\tilde{\mathcal{O}}_s(p; \alpha) = \mathcal{M}_{tot} \mathcal{B}_s(p; \alpha)$

$$\begin{aligned} \bar{\alpha}_\mu p^\mu \tilde{\mathcal{O}}_s(p; \alpha) &= 0 \\ \bar{\alpha}^\mu \bar{\alpha}_\mu \tilde{\mathcal{O}}_s(p; \alpha) &= 0 \\ \alpha^\mu \bar{\alpha}_\mu \tilde{\mathcal{O}}_s(p; \alpha) &= s \tilde{\mathcal{O}}_s(p; \alpha) \end{aligned}$$

→

$$\begin{aligned} (-\bar{\alpha}^0 p^0) \mathcal{B}_s(p; \alpha) &= 0 \\ (2\bar{\alpha}^1 \bar{\alpha}^2 - \bar{\alpha}^0 f(p; \bar{\alpha})) \mathcal{B}_s(p; \alpha) &= 0 \\ \alpha^\mu \bar{\alpha}_\mu \mathcal{B}_s(p; \alpha) &= s \mathcal{B}_s(p; \alpha) \end{aligned}$$

$$\begin{aligned} \tilde{\mathcal{O}}_s(p; \alpha) &= \mathcal{M}_{tot} \left( \mathcal{B}_s^{(1)}(p) (\alpha^1)^s + \mathcal{B}_s^{(2)}(p) (\alpha^2)^s \right) \\ &= \mathcal{B}_s^{(1)}(p) (\epsilon_\mu \alpha^\mu)^s + \mathcal{B}_s^{(2)}(p) (\epsilon_\mu^* \alpha^\mu)^s \end{aligned}$$

$$\epsilon(p) \equiv \frac{1}{\sqrt{2} |\vec{p}|} \left( \frac{\vec{p}^2}{\sqrt{-p^\mu p_\mu}}, \frac{p^0 p^1}{\sqrt{-p^\mu p_\mu}} + i p^2, \frac{p^0 p^2}{\sqrt{-p^\mu p_\mu}} - i p^1 \right) \quad : \text{polarization vector}$$



# Solution : spin-s primary operator

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□ Solution :  $\tilde{\mathcal{O}}_s(p; \alpha) = \mathcal{B}_s^{(1)}(p) (\epsilon_\mu \alpha^\mu)^s + \mathcal{B}_s^{(2)}(p) (\epsilon_\mu^* \alpha^\mu)^s$

■ Properties of polarization vector

$$\epsilon(p) \cdot \epsilon(p) = 0 \quad \epsilon(p) \cdot \epsilon(p) = 0 \quad p^\mu \cdot \epsilon_\mu(p) = 0$$

⇒  $\tilde{\mathcal{O}}_s(p; \alpha) = \tilde{\mathcal{O}}_s(p; \epsilon^*) (\epsilon_\mu \alpha^\mu)^s + \tilde{\mathcal{O}}_s(p; \epsilon) (\epsilon_\mu^* \alpha^\mu)^s$

□ Physical projected spin-s primary operator

□ Fourier transformation (  $p^0 > 0$  ,  $p_a^\mu = (|\vec{p}_a|, \vec{p}_a)$  )

$$\begin{aligned} \tilde{\mathcal{O}}_s(p; \alpha) &= 2p^0 \int d^3x e^{-ip \cdot x} \mathcal{O}^s(x; \alpha) \\ &= \int d\vec{p}_1 d\vec{p}_2 2\mathcal{J}(\vec{p}_1, \vec{p}_2) \delta^{(3)}(p_1^\mu + p_2^\mu - p^\mu) \frac{s!}{\Gamma(s + \frac{1}{2})} \\ &\quad \times (\alpha \cdot p_1 + \alpha \cdot p_2)^s P_s^{-\frac{1}{2}, -\frac{1}{2}} \left( \frac{\alpha \cdot p_1 - \alpha \cdot p_2}{\alpha \cdot p_1 + \alpha \cdot p_2} \right) a(\vec{p}_1, \vec{p}_2) \end{aligned}$$

: annihilation operator of bi-local field



# (Unconstrained) Primary spin-s Current

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## □ (Unconstrained) Primary spin-s operator

$$\begin{aligned}\tilde{\mathcal{O}}_s(p; \epsilon) = & \int d\vec{p}_1 d\vec{p}_2 \delta^{(3)}(p_1^\mu + p_2^\mu - p^\mu) \frac{(-1)^s 2^{\frac{3}{2}s-1}}{s!} \\ & \times (2 |\vec{p}_1| |\vec{p}_2| - 2\vec{p}_1 \cdot \vec{p}_2)^{\frac{s}{2}} e^{is\Theta(\vec{p}_1, \vec{p}_2)} a(p_1, p_2)\end{aligned}$$

$$\begin{aligned}\tilde{\mathcal{O}}_s(p; \epsilon^*) = & \int d\vec{p}_1 d\vec{p}_2 \delta^{(3)}(p_1^\mu + p_2^\mu - p^\mu) \frac{(-1)^s 2^{\frac{3}{2}s-1}}{s!} \\ & \times (2 |\vec{p}_1| |\vec{p}_2| - 2\vec{p}_1 \cdot \vec{p}_2)^{\frac{s}{2}} e^{-is\Theta(\vec{p}_1, \vec{p}_2)} a(p_1, p_2)\end{aligned}$$

$$\Theta(\vec{p}_1, \vec{p}_2) \equiv 2 \arctan \left( \frac{2\vec{p}_1 \times \vec{p}_2}{\sqrt{2\sqrt{\vec{p}_1^2 \vec{p}_2^2} - 2\vec{p}_1 \cdot \vec{p}_2} (\sqrt{\vec{p}_1^2} - \sqrt{\vec{p}_2^2})} \right)$$

: Equal to  $\theta$  map (up to sign)



# Other Dimension

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## □ $\text{AdS}_4 \times \text{S}^2 / \text{CFT}_3$ covariant gauge

[JY, A. Jevicki, J.P.Rodrigues, R.d.M.Koch, 2014, 1408.1255 ]:

- Representation of  $\text{SO}(2,3)$  for higher spin in the embedding space (5+5 dim) with constraints & Fronsdal Gauge
- Possible to find  $\text{SO}(2,3)$  generators in  $\text{AdS}_4 \times \text{S}^2$
- Too complicated to find map

## □ $\text{AdS}_3 \times \text{S}^1 / \text{CFT}_2$ covariant gauge :

- Representation of  $\text{SO}(2,2)$  for higher spin in the embedding space (4+4 dim) with constraints & Fronsdal Gauge
- Bi-local map

## □ $\text{AdS}_3 / \text{CFT}_2$ time-like :

- Bi-local map



# Thank You