

# FINITE FEATURES IN GRAVITY

## 5TH MANDELSTAM WORKSHOP

I MOTIVATION

12/1/2022  
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II FINITE BOUNDARIES IN  
LORENTZIAN GRAVITY

III FINITE EXAMPLES  
OF "AS/CFT."

2212.04944  
2209.06144  
2206.14146

VI EUCLIDIAN CONSIDERATIONS

discussions & collab.  
with:  
Galante, Harris, Hofman,  
Mühlmann, Steyer

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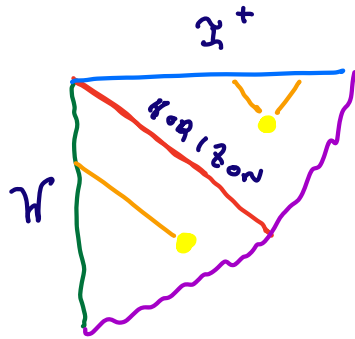
## MOTIVATION

SPACETIME MAY EXHIBIT NO ASYMPTOTIA

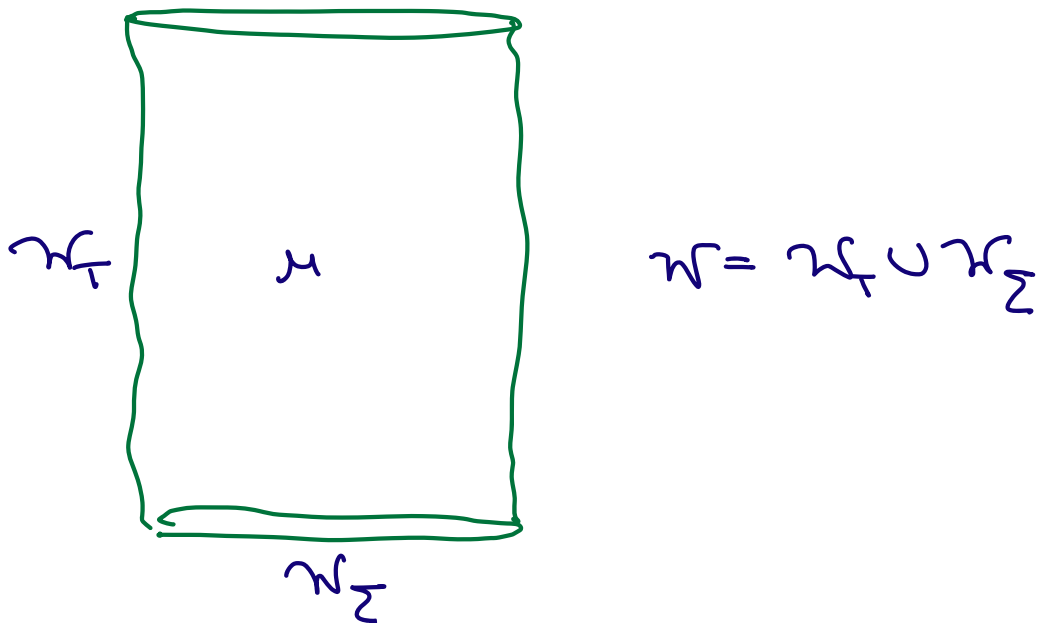
- e.g.
- \* COSMOLOGICAL SPACES WITH COMPACT CAUCHY SURFACES;
  - \* BIG BANG & BIG CRUNCH COSMOLOGIES
  - \* COSMOLOGICAL EVENT HORIZONS

IN SUCH A SITUATION A MORE  
QUASILocal DESCRIPTION MIGHT BE  
NECESSARY.

INDEED, IN OUR OWN WORLD, PROVIDED  
 $\Lambda > 0$  WE ARE CONFINED INSIDE THE  
DE SITTER HORIZON



FROM THE PERSPECTIVE OF GENERAL RELATIVITY  
 WE CAN ASK WHETHER A WELL-POSED  
 PROBLEM CAN BE FORMULATED ON A  
 MANIFOLD  $M$  WITH A TIMELIKE  
BOUNDARY  $\mathcal{W}$  OF FINITE SPATIAL SIZE.



DOES THERE EXIST DATA ON  $\mathcal{W}$   
 LEADING TO EXISTENCE AND/OR UNIQUENESS  
 OF SOLUTION IN  $M$

# DIRICHLET PROBLEM:

[Friedrich-Nagy;  
Anderson; An-Anderson;  
Witter;  
D.A., Galante, Nöhlmann]

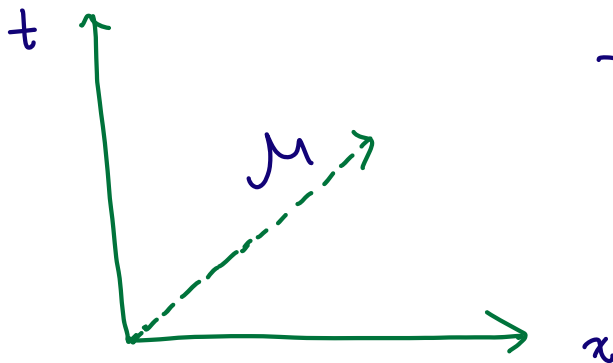
SPECIFY INDUCED METRIC  $h_{mn}$  ON  $\mathcal{W}_T$   
&  $(\tilde{h}_{\mu\nu}, \tilde{K}_{\mu\nu})$  ON  $\mathcal{W}_\Sigma$

THEOREM: (i) GENERIC DATA LEADS TO NON-EXISTENCE

[An-Anderson] (ii) DATA PERMITTING SOLUTION IN  $\mathcal{M}$   
PERMITS  $\infty$  MANY SOLUTIONS.

## SIMPLE EXAMPLE:

LINEARIZED GRAVITY NEAR MINKOWSKI



$$\mathcal{W}_T: \{x=0, t>0\}$$

$$\mathcal{W}_\Sigma: \{x>0, t=0\}$$

IN DE DONDER GAUGE:  $\nabla^\nu h_{\nu\mu} - \frac{1}{2} \partial_\mu h = 0$

& WITH DIRICHLET DATA:

$$h_{mn} = 0 \text{ ON } \mathcal{W}_T \quad \& \quad \tilde{h}_{\mu\nu} = \tilde{K}_{\mu\nu} = 0 \text{ ON } \mathcal{W}_\Sigma$$

ON FINDS  $\infty$  MANY SOLUTIONS!

MANY OF THESE CAN BE WRITTEN

AS: 
$$h_{\mu\nu} = \partial_\mu \xi_\nu + \partial_\nu \xi_\mu$$

HOWEVER:  $\xi_\mu$  IS NOT PURE GAUGE  
AS IT DOES NOT OBEY BOUNDARY  
CONDITIONS ON  $\mathcal{W}$ .

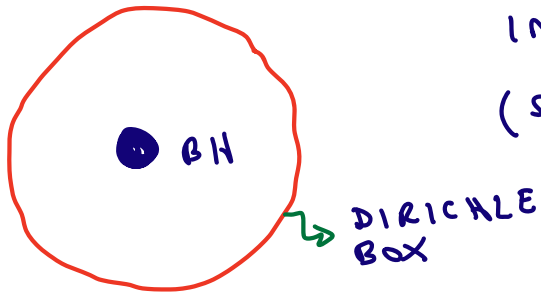
CONJECTURE: FIXING  $\tilde{h}$  &  $\text{tr} K$  ON  $\mathcal{N}_T$   
[An-Anderson] &  $(\tilde{h}_{\mu\nu}, \tilde{K}_{\mu\nu})$  ON  $\mathcal{N}_\Sigma$   
(WHERE  $\tilde{h}$  IS THE CONFORMAL CLASS OF  $h_{\mu\nu}$ )  
IS WELL-POSED.

- \* EVIDENCE PROVIDED AT LINEARIZED LEVEL
- \* REMINISCENT OF ADS BOUNDARY DATA  
BUT CRUCIALLY HERE  $\mathcal{W}$  IS A PRIORI OF  
FINITE SIZE

OPEN QUESTION: MATHEMATICAL STRUCTURE  
OF  $\mathcal{D}_{\text{dof}}$  OBSERVABLES  
(e.g. UNCLEAR IF AXIOMS OF LOCAL QFT APPLY)

# FINITE BOUNDARIES IN PHYSICS LITERATURE

York ~ 80s : BLACK HOLE THERMODYNAMICS  
IN A BOX.



(S-WAVE SECTOR / DIRICHLET)

't Hooft ; DAMOUR ; ... ;  
SUSSKIND-UGLUM-THORLACIOUS ;  
PRICE-THORNE ; ...

FLUID-GRAVITY LITERATURE

STRETCHED HORIZON  
PICTURE.



JACOBSON : EINSTEIN EQNS  
FROM LOCAL HORIZON  
THERMODYNAMICS

MORE RECENT ATTEMPTS TO APPLY TO COSMOLOGICAL  
HORIZONS :

[ Banks-Fischler ; ... ; D.A., Anous, Freedberg, Ng ; Banikashemi-Jacobson ;  
Coleman-Silvestein-Shyam-Torrobá ; Susskind ; ... ]

→ INCOMPRESSIBLE NAVIER-STOKES EQN

↕  
EINSTEIN-EQUATIONS

# HOLOGRAPHIC PERSPECTIVE

$AdS_{d+1} / CFT_d$   $\longleftrightarrow$  DUAL THEORY IS LOCAL FIELD THEORY.  
 $d \geq 2$ .

$\Rightarrow$  BULK MUST ACCOMMODATE SPATIAL CONTINUUM OF LOCAL OPERATORS

IT IS NATURAL TO POSTULATE THAT THESE OPERATORS ARE ACCOMMODATED AT ASYMPTOTIC BOUNDARY OF BULK SPACETIME

i.e. BULK THEORIES DUAL TO LOCAL QFTS SPATIAL SLICES OF INFINITE PROPER VOLUME; AT LEAST CLASSICALLY.

## QUANTUM MECHANICAL HOLOGRAPHY?

SOME HOLOGRAPHIC PAIRS INVOLVE QM DUAL

e.g.  $N$  DO BRANE QM  $\longleftrightarrow$  IIA DO BACKGROUND MATRIX QM  $\longleftrightarrow$  NON-CRITICAL STRING (near)-CFT,  $\longleftrightarrow$  (near)- $AdS_2$ .

IN THESE EXAMPLES BOUNDARY IS "WORLDLINE" & DUAL HAS FINITE D.O.F WHEN  $N < \infty$ .  
BUT STILL ASYMPTOTIC...

... NONETHELESS, IF ONE CAN ISOLATE SOME IR SUBSECTOR, ONE MIGHT HAVE A COMPLETELY "FINITE" HOLOGRAPHY.

SIMPLE  
EXAMPLE: RG FLOWS OF (near)-CFT<sub>1</sub>.

LET 
$$\hat{H}_g \equiv \sum_{i_1, \dots, i_g} \mathcal{J}_{i_1, \dots, i_g} \hat{\psi}_{i_1} \dots \hat{\psi}_{i_g}$$

WITH 
$$\langle \mathcal{J}_{i_1, \dots, i_g} \mathcal{J}_{j_1, \dots, j_g} \rangle = \frac{2^{g-1}}{g} \frac{g^2 (g-1)!}{N^{g-1}} \delta_{i_1 j_1} \dots \delta_{i_g j_g}$$
  
(à la Wigner)

AT ENERGY SCALES  $\ll g$

THEORY IS NEAR FIXED POINT WITH

$$\Delta_\psi \approx 1/g ;$$

$$\Theta^{(n)} \approx \psi_i \partial_t^n \psi_i \quad \text{with} \quad \Delta_n \gtrsim 1$$

Sachdev-Ye;  
Kitaev;  
Maldacena-Stanford;  
...

↳ SUBSET OF OPERATORS APPEARING IN  $\langle \psi_i \psi_j \psi_k \dots \rangle$

& SYMMETRY BREAKING SECTOR GOVERNED  
 BY SCHWARZIAN ACTIONS:

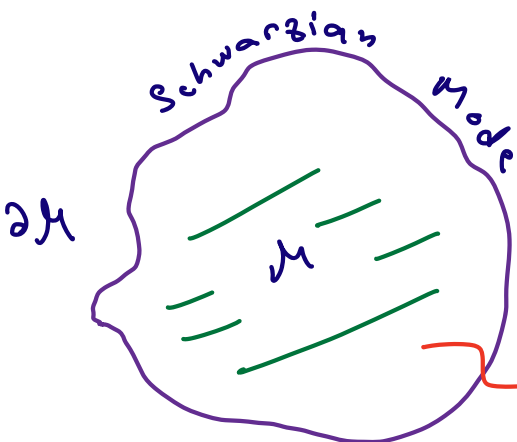
$$S_{Sch} = \# \frac{N}{J} \int_{\mathbb{R}} dt \text{ Sch} \{ f(t), t \}$$

CHARACTERISTIC SCALE OFFENDS CONFORMAL SYMMETRIES

LOW T THERMODYNAMICS FOLLOWS FROM  
 $\mathbb{R} \rightarrow S^1$  MAP

ONE FINDS:  $\lim_{N \rightarrow \infty} \frac{S}{N} = S_0 + \underbrace{c \cdot T}_{\text{SPECIFIC HEAT}} + \dots$   
 T=0 DEGENERACY

(SEMI) - HOLOGRAPHIC PICTURE



$$S = -\frac{1}{2\kappa} \int_{\partial M} d^2x \sqrt{g} \phi (R+2) - \frac{1}{\kappa} \int_{\partial M} d^2x \sqrt{h} \phi K$$

+ matter fields.

POINCARÉ DISK.

THE BULK GEOMETRY IS CLASSICALLY  
FEATURELESS.

STRUCTURE COMES FROM:

RELEVANT DEFORMATIONS

$\mathcal{O}^{(n)}$  ARE ALL IRRELEVANT!

HOWEVER  $\Delta \psi_i = 1/q$

→ AT LARGE  $q$  ROOM FOR MANY  
RELEVANT OPERATORS WITH  $\Delta < 1$ .

CONCRETELY WE STUDY:

[D.A., Galante, Sheeory]  
Jiang, Yang; ...

$$\hat{H} = \hat{H}_q + s \hat{H}_{\tilde{q}}$$

$$\tilde{q}/q < 1.$$

(NAÏVE) INTUITION IS THAT  $\hat{H}_{\tilde{q}} = \int_{i_1, \dots, i_{\tilde{q}}} \psi_{i_1} \dots \psi_{i_{\tilde{q}}}$

MIGHT FLOW TO RELEVANT DEFORMATIONS OF  
ORIGINAL (near)-CFT

WITH  $\Delta = \tilde{q}/q < 1$ .

RESULTS:

$\tilde{q}/q = 1/2$

PERMITS ANALYTIC TREATMENT  
FOR  $N \gg q \gg 1$ .

e.g.  $\frac{1}{N} \langle \psi_i(t) \psi_i(0) \rangle = \frac{sgt}{2} \left( \frac{1}{1 + g|t| \sqrt{1 + 4s^2} + s^2 g^2 |t|^2} \right)^{1/q}$

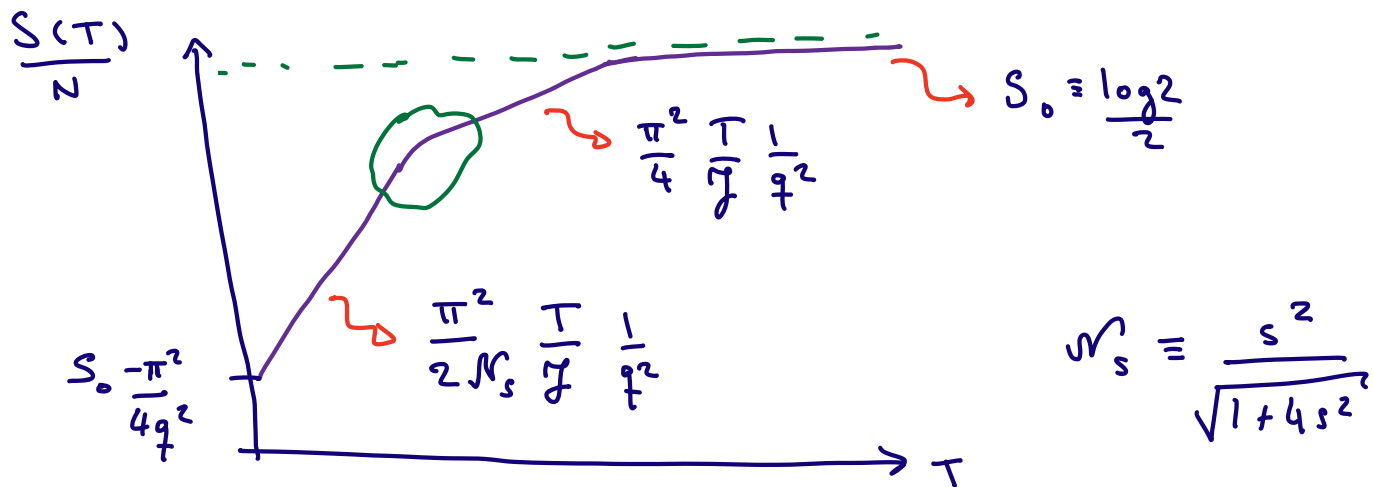
FOR SUFFICIENTLY SMALL S

A.  $\frac{1}{g \sqrt{1 + 4s^2}} \ll |t| \ll \frac{\sqrt{1 + 4s^2}}{s^2 g}$

B.  $|t| \gg \frac{\sqrt{1 + 4s^2}}{s^2 g}$

EXHIBIT  
NEAR-CFT  
BEHAVIOUR.

FINITE TEMPERATURE RG FLOW



TWO-REGIMES WITH LINEAR IN T ENTROPY

NUMERICAL EVIDENCE SUPPORTS

PICTURE FOR FINITE  $\tilde{g}$  & VARIOUS  $\tilde{g}/g < 1/2$

WHEN  $\tilde{g}/g < 2 \rightarrow$   $S_0$  depends on  $s$

MOREOVER, USING:

CONFORMAL PERTURBATION THEORY

WE CAN STUDY CORRECTION TO THERMODYNAMICS  
NEAR FIXED POINT A.

CORRECTION (NON-ANALYTIC BEHAVIOUR IN  $T$ )

COMPATIBLE WITH:

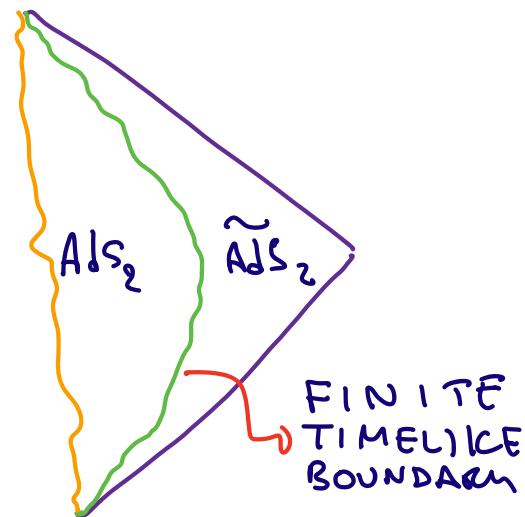
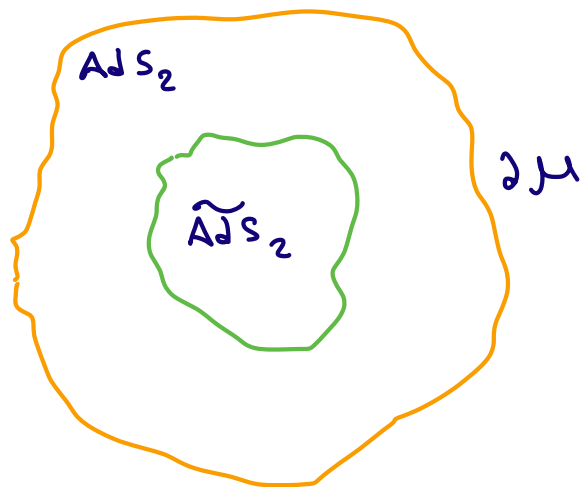
$$\Delta_{\text{rel}} = 1/n$$

$\rightarrow$  PARTONIC  
PICTURE

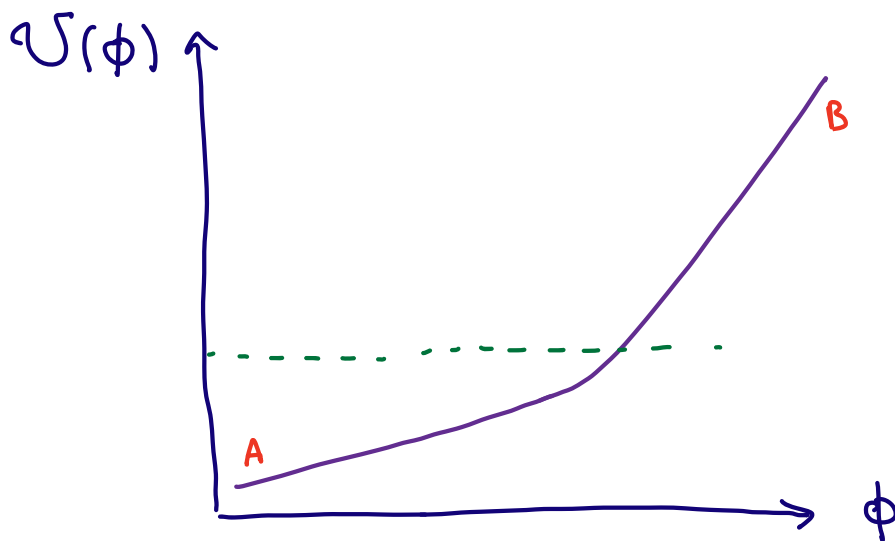
i.e.  $(\gamma/T)^{2-2/n}$

ALSO FIXED POINT B HAS EMERGENT  
SCHWARZIAN MODE EMBEDDED ENTIRELY  
INSIDE STRONG COUPLING SECTOR.

# GEOMETRIZATION OF RG FLOW



$$S_{\text{deformed}} = -\frac{1}{2\kappa} \int d^2x \sqrt{g} [\phi R + \mathcal{V}(\phi)] - \frac{1}{\kappa} \int dx \sqrt{h} \phi K$$



$$\boxed{R = -\partial_\phi \mathcal{V}(\phi)} \quad \rightarrow \quad 2 \text{ AdS}_2 \text{ REGIONS}$$

BOUNDARY OF INTERIOR  $\text{AdS}_2$  GOVERNED BY EMERGENT SCHWARZIAN MODE RESIDING ENTIRELY WITHIN GEOMETRIC DESCRIPTION.

IN LORENTZIAN SIGNATURE WE HAVE  
A FINITE TIMELIKE BOUNDARY

APPENDED WITH SIMPLE OBSERVABLES

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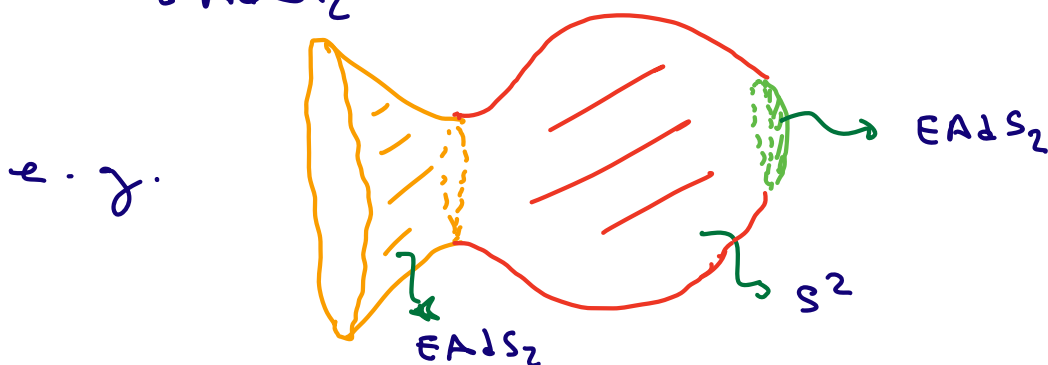
LANDSCAPE OF RG FLOWS

IN A SIMILAR SPIRIT ONE CAN  
IMAGINE RICHER CLASSES OF RG-FLOWS

e.g. 
$$\hat{H} = \sum_{i=1}^n s_i \hat{H}_{f_i}$$
 with  $f_1 > f_2 > \dots > f_n$ .

LEADING TO INCREASINGLY RICH STRUCTURE  
SUSPENDED IN INTERIOR OF  $AdS_2$

OF PARTICULAR INTEREST IS TO OBTAIN  
A PORTION OF  $S^2$  SUSPENDED BETWEEN  
TWO  $EAdS_2$  GEOMETRIES



THESE GEOMETRIES IN LORENTZIAN SIGNATURE  
INTERPOLATE BETWEEN  $AdS_2$  &  $dS_2$

[D.A., Hofman; Galante; Harris]

SOME MICROPHYSICAL EVIDENCE FOR THESE  
COMES FROM SIMPLIFIED DESCRIPTION OF  
JT GRAVITY IN TERMS OF MATRIX  
INTEGRALS WITH GENERALIZED POTENTIAL.

[Witten; Maxfield-Turiaci]

NOT YET EMBEDDED IN SYK...

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## GENERAL COMMENTS ON SPHERE

MOST "FINITE" GEOMETRY SUPPORTED BY  
GRAVITY. SADDLE POINT SOLUTION TO

$$\mathcal{Z} = \int [Dg] e^{-\frac{1}{16\pi G} \int d^4x \sqrt{g} (-R + 2\Lambda)}$$

$\Lambda > 0$  CRUCIAL TO SUPPRESS LARGE  
VOLUME. Diff, Gauge, Field Redef. invariant.

PHYSICAL MEANING OBSCURED BY  
ABSENCE OF BOUNDARY.

GIBBONS-HAWKING POSTULATE

$$\tilde{Z} = e^S$$

ENTROPY OF DE SITTER  
HORIZON.

CLASSICALLY

$$S_0 = \frac{A_{hor}}{4G} = \frac{3\pi}{\Lambda G}$$

\* ABSENCE OF BOUNDARY  $\Rightarrow$  NO TEMPERATURE  
PARAMETER.

\* ALSO ADM ENERGY VANISHES.

SOME SIMILARITIES TO BPS STORY OF  
BLACK HOLES

$$S \stackrel{D}{=} \dim \mathcal{R}_{dS}$$

[Fischler - Banks - Fiol, ...]  
Verlinde, Susskind, ...

GENERAL QUANTUM STRUCTURE:

$$\log \tilde{Z} = S_0 - \underbrace{\frac{\dim \text{SO}(5)}{2}}_{\text{VOLUME OF RESIDUAL}} \log S_0 - \left( \frac{331}{90} + \frac{8}{3} \right) \log \frac{3}{\Lambda_{\text{eff}}^2} + c + \mathcal{O}(S_0^{-1})$$

ONE-LOOP PIECE: [D.A., Denef, Law, Sun]

$$S^{(1)} = \int_0^{\infty} d\omega \left( -\log \left[ e^{b\omega/2} - e^{-b\omega/2} \right] \rho(\omega) \right)$$

WHERE  $\rho(\omega)$  FOLLOWS FROM HARIHAR-CHANDRA  
GROUP CHARACTER OF  $SO(D, 1)$ ;  $b = 2\pi l$

→ APPEARANCE OF THERMAL STRUCTURE

MAKING SENSE OF  $\mathbb{Z}$  (THROUGH A  
VARIETY OF TOOLS) IS AN ADDITIONAL  
PATH TOWARD FINITE HOLOGRAPHY

e.g.

IN 3D:  $T\bar{T}$  DEFORMATIONS OF  $ADS_3 / CFT_2$   
INTO  $DS_3$  PROPOSED BY  
[Coleman-Mazenc-Shyam-Silverstein-Soni-Torrobá-Yang]  
REPRODUCE LOGARITHMIC CORRECTIONS.

IN 2D: CONNECTION TO OLD MATRIX  
INTEGRALS ...

[Polyakov-Knizhnik-Zamolodchikov-Kazakov - ... - D.A., Mühlmann]  
Ambjörn, Distler-Kawai, David

# OUTLOOK

\* NEW AVENUES (GR; HOLOGRAPHY;  
EUCLIDEAN GRAVITY, ...)

TO EXPLORE FINITE FEATURES IN Q.G.

\* MUCH TO EXPLORE, CONNECTING DIFFERENT PICTURES.

\* ULTIMATELY, IT MAY BE ABOUT A CLEARER UNDERSTANDING OF THE ROLE OF AN OBSERVER IN Q.G. ?

THANK YOU!